

Multicritical dissipative phase transitions in the anisotropic open quantum Rabi model



Guitao Lyu¹, Korbinian Kottmann², Martin B. Plenio³, and Myung-Joong Hwang¹

¹ Division of Natural and Applied Sciences, Duke Kunshan University, China

² ICFO - Institut de Ciències Fòniques, The Barcelona Institute of Science and Technology, Spain

³ Institute of Theoretical Physics and IQST, Ulm University, Germany

Abstract

We study a generalized open quantum Rabi model with varying coupling strengths in the rotating and counter-rotating terms. Using both semiclassical and quantum approaches, we find a rich phase diagram resulting from the interplay between the anisotropic interaction and the dissipation.

First, there exists a bistable phase where both the normal and superradiant phases are stable. Second, there are multicritical points where the phase boundaries for the first- and second-order phase transitions meet. We show that a new set of critical exponents governs the scaling of the multicritical points, which could be utilized for developing critical sensing protocols.

Model: anisotropic open QRM

$$H = \omega_0 a^\dagger a + \frac{\Omega}{2} \sigma_z - \lambda_r (a \sigma_+ + a^\dagger \sigma_-) - \lambda_{cr} (a \sigma_- + a^\dagger \sigma_+)$$

$$\dot{\rho} = -i[H, \rho] + \kappa(2a\rho a^\dagger - a^\dagger a \rho - \rho a^\dagger a)$$

Where the anisotropic coupling strengths are denoted by λ_r and λ_{cr} for the rotating and counterrotating terms, and κ is the damping rate of the harmonic oscillator.

We consider that $\Omega/\omega_0 \rightarrow \infty$, which serves as the thermodynamical limit. In this limit, the (open) Rabi model exhibits a (dissipative) superradiant phase transition, as predicted in PRL2015 and PRA2018 by Hwang *et al.*

Semiclassical analysis

- We obtain a mean-field solution for the steady states by solving the semiclassical equations of motion:

$$\begin{aligned} \langle \dot{a} \rangle &= -i(\omega_0 - i\kappa) \langle a \rangle + i(\lambda_{cr} \langle \sigma_+ \rangle + \lambda_r \langle \sigma_- \rangle), \\ \langle \dot{\sigma}_+ \rangle &= i\Omega \langle \sigma_+ \rangle + i(\lambda_{cr} \langle a \rangle + \lambda_r \langle a^\dagger \rangle) \langle \sigma_z \rangle, \\ \langle \dot{\sigma}_z \rangle &= i2(\lambda_r \langle a \rangle + \lambda_{cr} \langle a^\dagger \rangle) \langle \sigma_+ \rangle + \text{c.c.} \end{aligned} \quad \begin{aligned} \langle \dot{a} \rangle &\equiv \frac{d}{dt} \langle a \rangle \\ \langle a \rangle &\equiv \text{Tr}[\rho a] \end{aligned}$$

The solution for the **normal phase (NP)**:

$$s_z = 1, x = y = s_x = s_y = 0$$

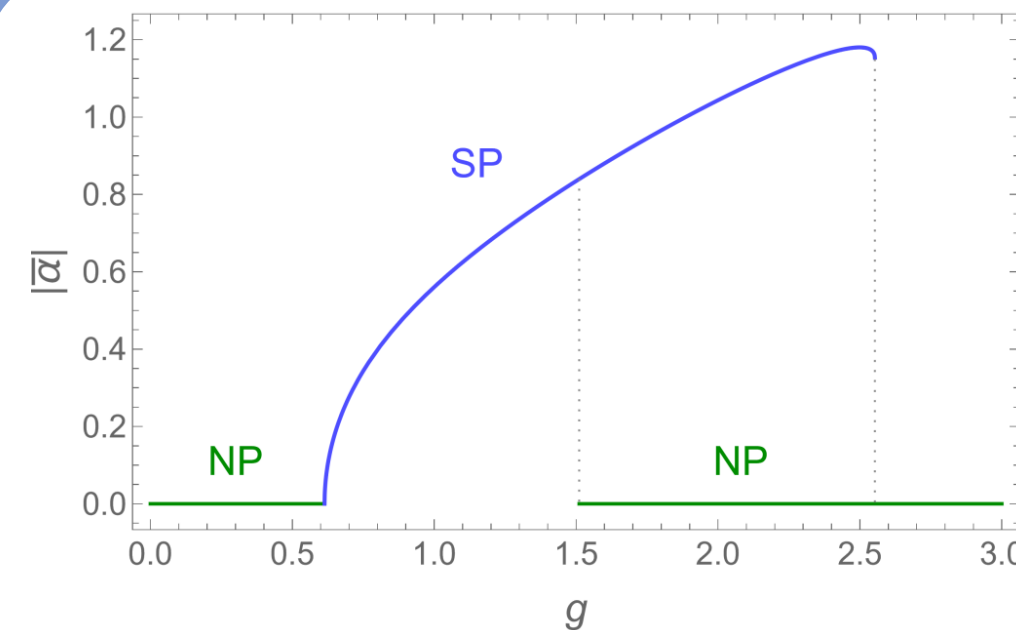
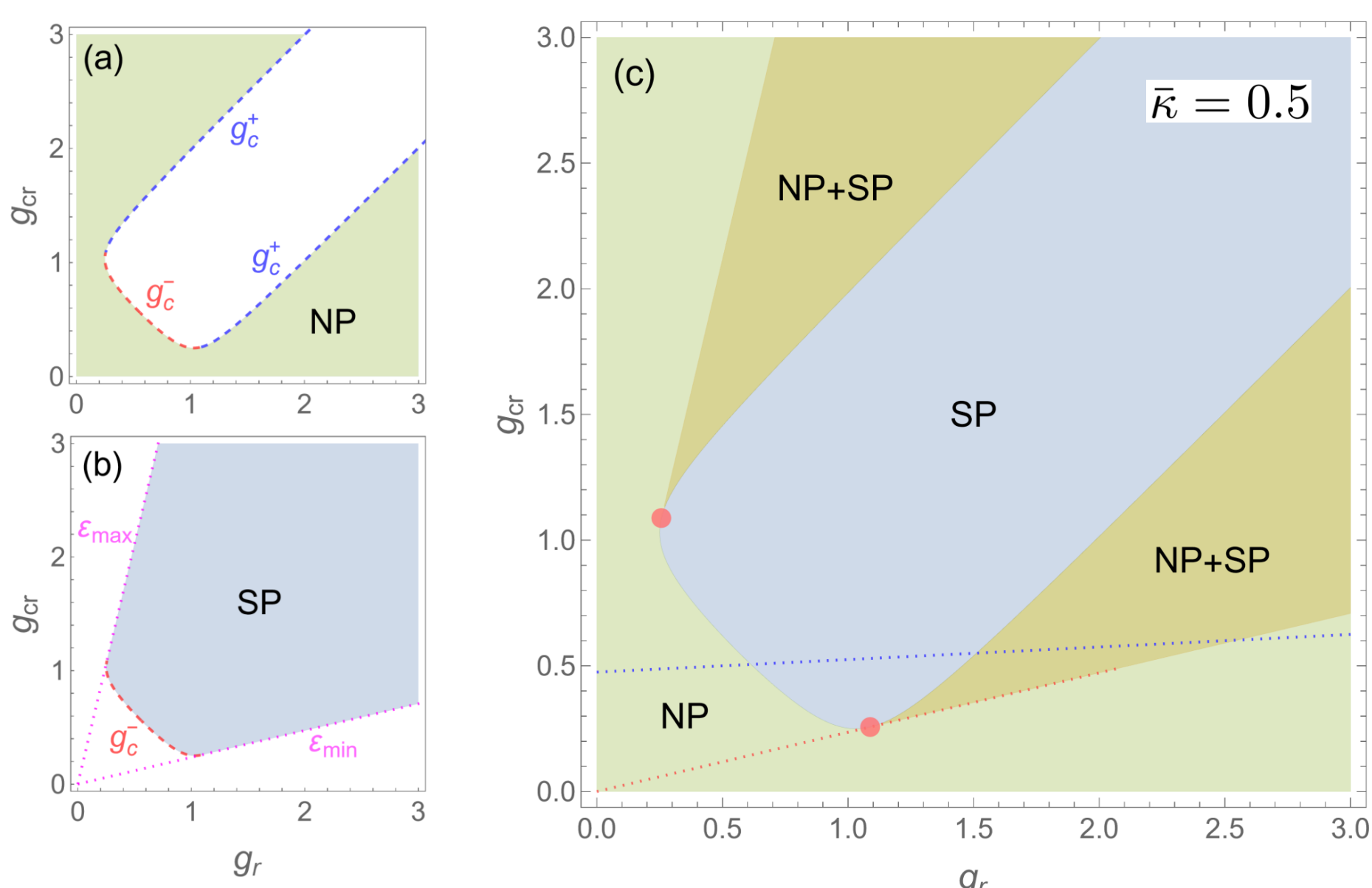
The solution for the **superradiant phase (SP)**:

$$s_z = -\frac{(1 + \varepsilon^2) \pm \sqrt{4\varepsilon^2 - \bar{\kappa}^2(1 - \varepsilon^2)^2}}{(1 - \varepsilon^2)^2 g^2}$$

$$s_{x,y,z} \equiv \langle \sigma_{x,y,z} \rangle,$$

$$\begin{aligned} g_r &\equiv \lambda_r / \sqrt{\omega_0 \Omega}, \\ g_{cr} &\equiv \lambda_{cr} / \sqrt{\omega_0 \Omega}, \\ \bar{\kappa} &\equiv \kappa / \omega_0, \\ g_r &\equiv g, \quad g_{cr} \equiv \varepsilon g \end{aligned}$$

- After a stability analysis for above mean-field solutions, we obtain the phase diagram of the system. (The following quantum solutions agree with the phase diagram.)



Dissipative phase transitions. The absolute value of a as g going along the blue dotted line in Fig. 1(c).

Quantum fluctuations

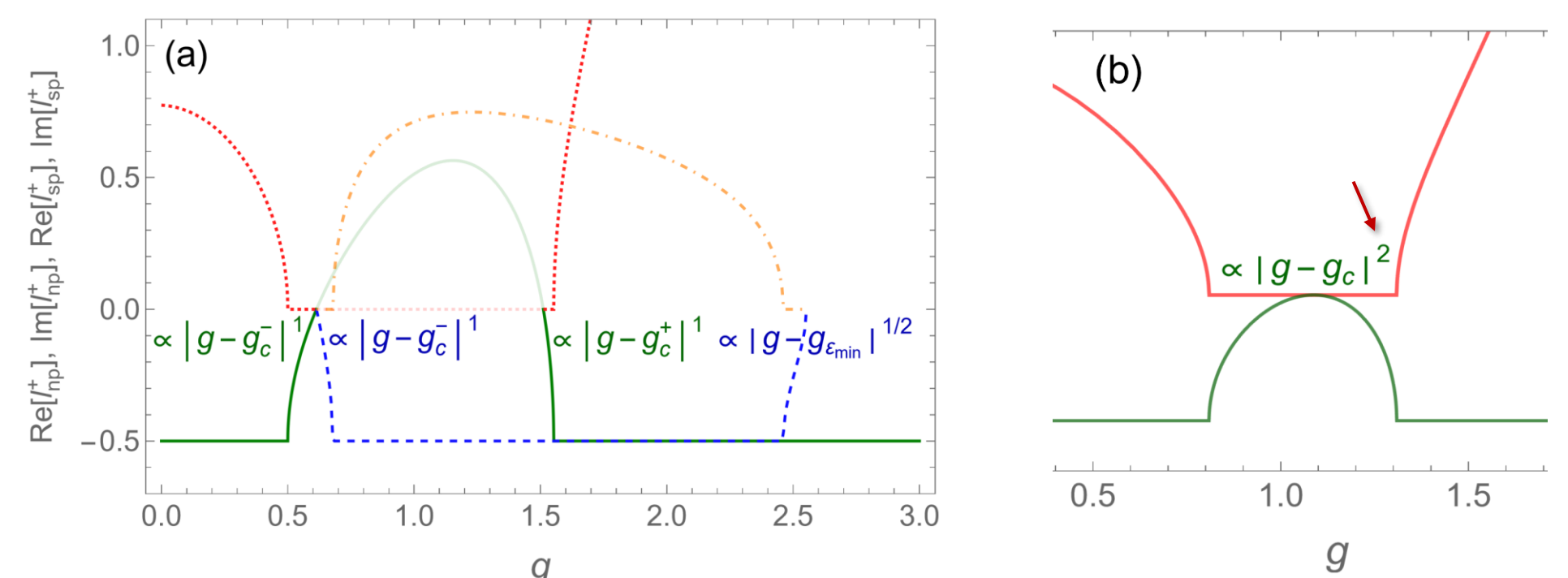
By performing the Schrieffer-Wolff (SW) transformation to the Hamiltonian and projecting it to the spin-down subspace, we obtain the effective Hamiltonian for the NP. (Similarly for the SP, but a displacement to the master equation is necessary.)

$$H_{np} = \omega_0 a^\dagger a - \omega_0 [g_r^2 a^\dagger a + g_{cr}^2 a a^\dagger + g_r g_{cr} (a^{\dagger 2} + a^2)]$$

- First moment of the oscillator operators:

$$\dot{\mathbf{a}} = L_{np} \mathbf{a} \quad \mathbf{a} = (\langle a \rangle, \langle a^\dagger \rangle)^\top$$

$$l_{np}^\pm = -\bar{\kappa} \pm \sqrt{2g^2(1 + \varepsilon^2) - g^4(1 - \varepsilon^2)^2 - 1}$$

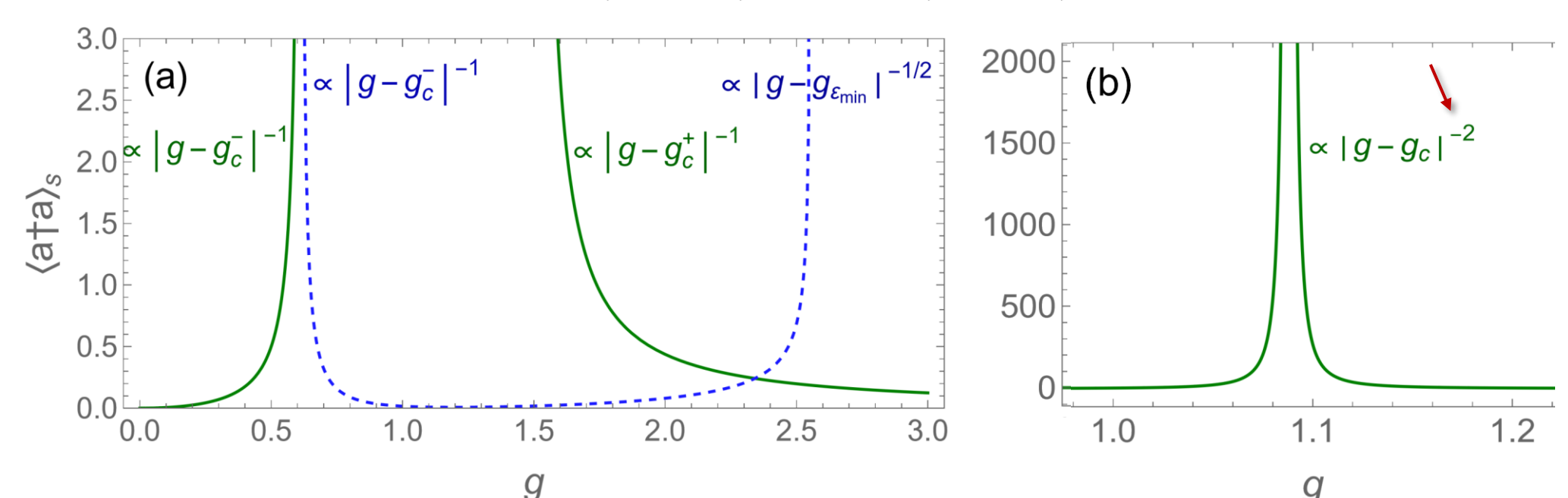


Asymptotic decay rate of the NP (solid line) and the SP (dashed line). (a) [(b)] is for the case g going along the blue [red] dotted line in Fig. 1(c).

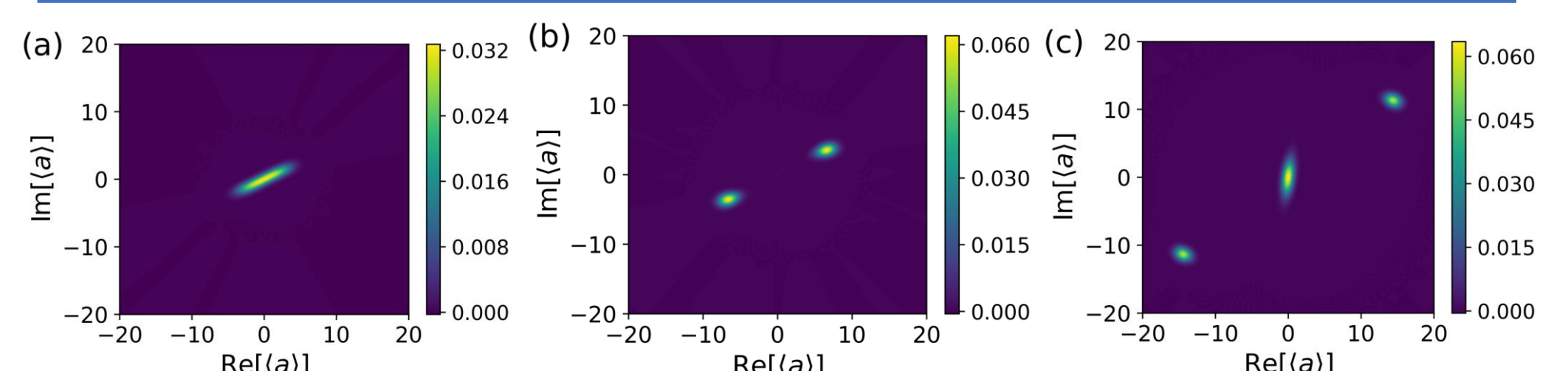
- Second moment of the oscillator operators:

$$\dot{\mathbf{s}} = M_{np} \mathbf{s} + Y_{np} \quad \mathbf{s} = (\langle a^\dagger a \rangle, \langle a^2 \rangle, \langle a^{\dagger 2} \rangle)^\top$$

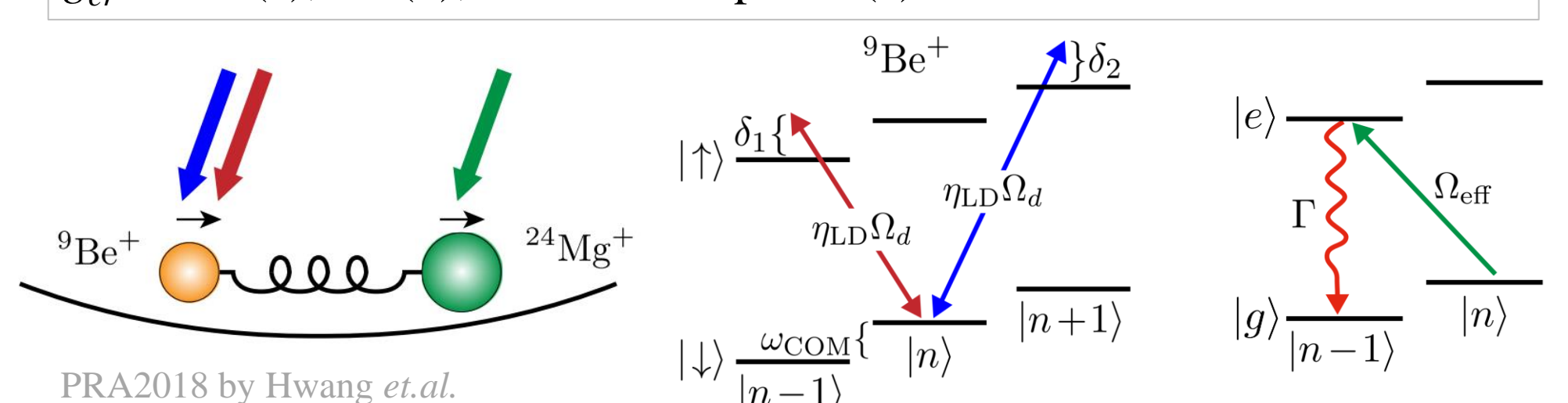
$$\langle a^\dagger a \rangle_s = \frac{2g^4 \varepsilon^2}{g^4(1 - \varepsilon^2)^2 - 2g^2(1 + \varepsilon^2) + \bar{\kappa}^2 + 1}$$



The corresponding Scaling of the oscillator excitation number.



Numerical simulation of the master equation for finite Ω/ω_0 . Wigner functions of the cavity field of the steady state for various g , and g_{cr} in NP (a), SP (b), and bistable phase (c).



PRA2018 by Hwang *et al.*